

# **10 - Natural (Free) Convection**



This is called free or natural convection. The solution is quite complicated.

When  $T_w - T_{fluid}$  is not so large, we may use equations for constant property, incompressible fluid with bouyancy forces when the density change is included in the body force term in the momentum equation.



### **ME – 212 THERMO-FLUIDS ENGINEERING II**



Buoyancy forces are responsible for the fluid motion in natural convection.

Viscous forces oppose the fluid motion.



Buoyancy forces are expressed in terms of fluid temperature differences through the volume expansion coefficient

$$\beta = \frac{1}{V} \left. \frac{\partial V}{\partial T} \right|_{P} = \frac{-1}{\rho} \left. \frac{\partial \rho}{\partial T} \right|_{P}$$



Volume Expansion Coefficient:

$$\beta = \frac{1}{V} \left. \frac{\partial V}{\partial T} \right|_{P} = \frac{-1}{\rho} \left. \frac{\partial \rho}{\partial T} \right|_{P}$$

Т

The volume expansion coefficient can be expressed approximately by replacing differential quantities by differences as

$$\beta \approx \frac{-1}{\rho} \frac{\Delta \rho}{\Delta T} = \frac{-1}{\rho} \left( \frac{\rho_{\infty} - \rho}{T_{\infty} - T} \right) \quad \text{at constant P}$$

$$\rho_{\infty} - \rho = \rho \beta (T_{\infty} - T) \quad \text{at constant P}$$
For an Ideal gas, (P V = m R T) 
$$\beta_{ideal gas} = \frac{1}{T} \quad \text{in } K^{-1}$$





# **10.1 Boundary Layer Equations**

Consider a vertical hot flat plate immersed in a quiescent fluid body.

Assumptions:

- steady,

– laminar,

- two-dimensional,
- Newtonian fluid, and
- constant properties, except the density difference (ρ ρ<sub>∞</sub>) (Boussinesq approximation).



#### **ME – 212 THERMO-FLUIDS ENGINEERING II**



Joseph Valentin Boussinesq

# French mathematician and physicist

1842 - 1929











Conservation of momentum  $\rho \left( u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} \right) = \mu \frac{\partial^2 u}{\partial y^2} - \frac{\partial P}{\partial x} - \rho g$ in the x direction in the boundary layer x momentum equation in the quiescent  $\frac{\partial P}{\partial x} = -\rho_{\infty} g$ fluid outside the boundary layer (u = 0) $\frac{\partial \mathbf{v}}{\partial \mathbf{x}} \cong \frac{\partial \mathbf{v}}{\partial \mathbf{v}} \cong \mathbf{0}$ In the boundary layer, v << u and  $\frac{\partial P}{\partial v} = 0 \qquad \Rightarrow \qquad \frac{\partial P}{\partial x} = \frac{\partial P_{\infty}}{\partial x} = -\rho_{\infty} g$ Force balance in the y direction: Substitute:  $\rho \left( u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial v} \right) = \mu \frac{\partial^2 u}{\partial v^2} + (\rho_{\infty} - \rho) g$ 



x momentum equation :

$$\rho \left( \mathsf{u} \; \frac{\partial \mathsf{u}}{\partial \mathsf{x}} + \mathsf{v} \; \frac{\partial \mathsf{u}}{\partial \mathsf{y}} \right) = \mu \; \frac{\partial^2 \mathsf{u}}{\partial \mathsf{y}^2} + \left( \rho_{\infty} - \rho \right) \, \mathsf{g}$$

Substitute:  $\rho_{\infty} - \rho = \rho \beta (T_{\infty} - T)$ 

$$\left(u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y}\right) = v \frac{\partial^2 u}{\partial y^2} + g \beta (T - T_{\infty})$$

x momentum equation in non-dimensional form, when:

$$\mathsf{T}^+ = \frac{\mathsf{T} - \mathsf{T}_{\infty}}{\mathsf{T}_{\mathrm{s}} - \mathsf{T}_{\infty}}$$

$$\left(u^{+} \frac{\partial u^{+}}{\partial x^{+}} + v^{+} \frac{\partial u^{+}}{\partial y^{+}}\right) = \frac{1}{\operatorname{Re}_{L}} \frac{\partial^{2} u^{+}}{\partial (y^{+})^{2}} + \underbrace{\frac{g \beta (T_{s} - T_{s}) L_{c}^{3}}{v^{2}} \frac{T^{+}}{\operatorname{Re}_{L}^{2}}}_{V^{2}}$$



The new non-dimensional parameter represents the natural convection effects, and is called the Grashof number, Gr.



$$\mathsf{Gr}_{L} = \frac{\mathsf{g}\,\beta\,(\mathsf{T}_{s}\,\mathsf{-}\,\mathsf{T}_{\infty})\,\mathsf{L}_{c}^{3}}{\nu^{2}}$$

$$Gr_{L} = \frac{Buoyancy Force}{Viscous Force}$$

The flow regime in natural convection is governed by the *Grashof number* 

 $Gr_L > 10^9$  Flow is turbulent





Franz Grashof German Engineer 1826 – 1893





If 
$$\frac{Gr_L}{Re^2} \ll 1$$
 No free convection, no buoyancy

If 
$$\frac{\text{Gr}_{L}}{\text{Re}^{2}} >> 1$$
 Buoyancy dominates



### **Natural Convection over Surfaces**

In free (or natural) convection, buoyancy forces are the only driving force that generates the flow field because there is no external field. Therefore, Re cannot be an independent parameter. Instead, Gr is.

Nu = f(Gr, Pr)

Define **Rayleigh number**: Ra = Gr Pr =  $\frac{g \beta (T_s - T_{\infty}) L_c^3}{v \alpha}$ 

Sometimes, Ra is used in the relations: Nu = f(Ra)

See Table 12-3 for the correlations over various surfaces.





John William Strutt 3<sup>rd</sup> Baron Rayleigh British Mathematician and Physicist 1842 – 1919



# Example 1



Properties of water

Cylinder in water		Evaluate properties of water
D = 2.5 cm	h = ?	at the mean temperature
L = 0.5 cm	Q = ?	$\overline{T} = \frac{T_w + T_{\infty}}{T_w + T_{\infty}}$
$T_{\infty} = 20 \ ^{\circ}C$		2 55 + 20
$T_w = 55 \ ^\circ C$		$=\frac{33720}{2}=37.5$ °C
c <sub>p</sub> = 4.174 kJ/kg.k	K	k = 0.63 W/m.K
- ρ = 993 kg/m³		Pr = 4.53

at 37.5 °C

$$\rho = 993 \text{ kg/m}^3 \qquad \text{Pr} = 4.53$$

$$\mu = 6.82 \ 10^{-4} \text{ kg/m.s} \qquad \frac{\text{g } \beta \ \rho^2 \ \text{c}_p}{\mu \ \text{k}} = 3.3 \ 10^{10} \ \text{m}^{-3}.\text{K}^{-1}$$



$$Gr_{L} Pr = Ra_{L} = \frac{g \beta \rho^{2} c_{p}}{\mu k} L^{3} (T_{w} - T_{\infty})$$

 $Ra_{L} = (3.3 \ 10^{10}) \ (0.5)^{3} \ (55 - 20) = 1.03 \ 10^{11} > 10^{9} => Turbulent$ 

$$\frac{g \beta \rho^{2} c_{p}}{\mu k} = 3.3 \ 10^{10} \text{ m}^{-3} \text{.K}^{-1}$$

$$\beta = 3.5 \ 10^{-4} \text{ K}^{-1} \text{ or } ^{\circ}\text{C}^{-1}$$

$$g = 9.8 \text{ m/s}$$

 $Nu_{mean} = \frac{h_{mean} L}{k} = C (Gr_L Pr)^m$  Use Table => C = 0.1 and m = 1/3

Nu<sub>m</sub> = (0.1) 
$$(1.03 \ 10^{11})^{1/3}$$
 = 46.875  $h_m = \frac{0.63}{0.5} (46.875) = 59.06 \ W/m^2.K$ 

 $Q = h_m A (T_w - T_\infty) = (59.06) \pi (0.025) (0.5) (55 - 20) = 58.51 W$ 



## Example 2



Evaluate properties of air at the mean temperature  $\overline{T} = \frac{T_w + T_\infty}{2} = \frac{65 + 10}{2}$ 

$$c_p = 1.006 \text{ kJ/kg.K}$$
 $k = 0.027 \text{ W/m.K}$ Properties of air $\rho = 1.18 \text{ kg/m}^3$  $Pr = 0.707$ at 310.5 K $\mu = 2.0 \ 10^{-5} \text{ kg/m.s}$  $\alpha = 0.23 \ 10^{-4} \text{ m}^2/\text{s}$  $v = 15.7 \ 10^{-6} \text{ m}2/\text{s}$ 



$$Gr_{L} = \frac{g \beta (T_{w} - T_{\infty}) L^{3}}{\upsilon^{2}} = \frac{(9.8) \frac{1}{283} (65 - 10) (0.3^{3})}{(15.7 \ 10^{-6})^{2}} = 2.086 \ 10^{8}$$

 $Gr_{L} Pr = (2.086 \ 10^{8}) \ (0.707) = 1.475 \ 10^{8} < 10^{9}$  Flow is Laminar

Use Table  
for Ra 
$$\approx 1.5 \ 10^8$$
  $\begin{bmatrix} C = 0.59 \\ m = 1/4 \end{bmatrix}$   $\overline{Nu} = \frac{\overline{h} L}{k} = (0.59) \ (Ra)^{1/4} = (0.59) \ (1.5 \ 10^8)^{1/4} = 65.02$ 

$$\overline{h} = \frac{k}{L} \ \overline{N}u = \frac{0.027}{0.3} \ (65.02) = 5.85 \ W/m^2.K$$

Alternatively, we may use Table 7.2 to find h.



11 – Boiling and Condensation

